

Yielding Behavior in Injectable Hydrogels from Telechelic Proteins

Bradley D. Olsen, Julia A. Kornfield,* and David A. Tirrell*

Division of Chemistry and Chemical Engineering, California Institute of Technology, Pasadena, California 91125, United States

Received June 28, 2010; Revised Manuscript Received September 27, 2010

ABSTRACT: Injectable hydrogels show substantial promise for use in minimally invasive tissue engineering and drug delivery procedures. A new injectable hydrogel material, developed from recombinant telechelic proteins expressed in *E. coli*, demonstrates shear thinning by 3 orders of magnitude at large strains. Large-amplitude oscillatory shear illustrates that shear thinning is due to yielding within the bulk of the gel, and the rheological response and flow profiles are consistent with a shear-banding mechanism for yielding. The sharp yielding transition and large magnitude of the apparent shear thinning allow gels to be injected through narrow gauge needles with only gentle hand pressure. After injection the gels reset to full elastic strength in seconds due to rapid re-formation of the physical network junctions, allowing self-supporting structures to be formed. The shear thinning and recovery behavior is largely independent of the midblock length, enabling genetic engineering to be used to control the equilibrium modulus of the gel without loss of the characteristic yielding behavior. The shear-banding mechanism localizes deformation during flow into narrow regions of the gels, allowing more than 95% of seeded cells to survive the injection process.

Introduction

Hydrogels are widely used in tissue engineering because of the similarity between their properties and those of native tissue.³ Injectable hydrogels offer added advantages, in that they can be used in minimally invasive surgical procedures.^{1,2} In addition to meeting traditional tissue engineering design criteria (including tunable gel modulus and biological functionality^{3–5}), injectable gels must flow under modest pressure and set rapidly at the target site. A variety of injectable materials have been developed either by chemically cross-linking the gel after injection^{6–8} or by engineering the gel to undergo a temperature-,^{9–11} pH-,¹² or ion-sensitive² transition between the sol and gel states. Alternatively, a few injectable materials based on shear-thinning hydrogels have been reported,^{12–19} where a drop in the modulus of the gel at high shear rates allows it to flow during injection. Recovery of the elastic modulus after shear is potentially more rapid than chemical cross-linking or physically responsive gelation, providing for more uniform cell distributions in the injected gel¹⁵ and greater control over placement of the material.¹² In addition, shear thinning materials can be engineered to form stable gel phases across wide ranges of temperatures and solution conditions, allowing environmentally stable formulations to be prepared long before use.

Although the biological properties and linear mechanics of tissue engineering hydrogels have been extensively investigated, the nonlinear rheology that leads to shear thinning and injectability is less well understood. Dissociation of the physically associating groups under shear is clearly implicated in shear thinning behavior,^{12–16} however, the flow profiles and mechanisms of thinning are often unknown. In particular, it has been extensively demonstrated for a wide variety of entangled systems, in particular polymer melts and wormlike micellar solutions, that inhomogeneous flow profiles may occur due to instabilities in the stress–strain curve.^{20–25} These instabilities are manifest as shear bands, areas of high deformation rate where the entanglements

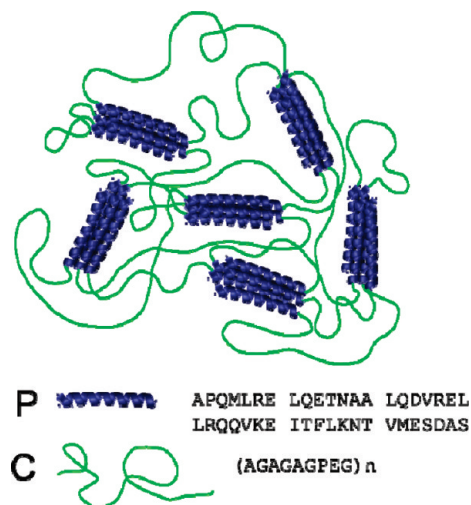


Figure 1. Structure of telechelic protein hydrogels. Telechelic proteins with coiled-coil end blocks and flexible polyelectrolyte midblocks associate in solution to form networks with junctions defined by the coiled-coil association state and connected by monodisperse linkers.

are preferentially disrupted, that accommodate the majority of the strain applied to the system.

Tirrell and co-workers reported a system of biofunctional physical hydrogels designed from telechelic proteins²⁶ in which coiled-coil end blocks are joined by a flexible polyanionic linker (Figure 1), and the coiled-coil motif has been exploited in both telechelic^{27–30} and graft polymers^{31–33} to produce a wide variety of protein hydrogels. Protein engineering allows control of the coiled-coil association multiplicity in the network junctions, and changes in the coiled-coil sequence^{29,30} and midblock molecular weight^{28,34} can be used to vary the modulus of the hydrogel through changes in the relative fractions of looped and bridged chains in the gel. Functional sequences such as the RGD cell binding domain can be incorporated into the midblock through

*To whom correspondence should be addressed. E-mail: tirrell@caltech.edu (D.A.T.) or jak@cheme.caltech.edu (J.A.K.).

genetic engineering without affecting the physical properties of the materials.³⁵ Gel structures based on similar design paradigms have been reported for block copolypeptides,³⁶ elastin-mimetic triblock copolymers,^{37–39} and telechelic proteins that use collagen-like blocks to form associating helical domains.^{40,41}

Here we evaluate a model system of telechelic protein hydrogels as injectable materials, demonstrating striking shear-thinning behavior and nearly instantaneous recovery of full elastic strength after cessation of shear. The observed flow behavior is explained in terms of the formation of shear bands in the gel at high stress. Yielding of the gel promotes rapid recovery of the elastic modulus and allows encapsulated cells to survive injection by distributing strain inhomogeneously in the flowing material.

Experimental Methods

For expression of the proteins PC₁₀P and PC₃₀P, the corresponding genes were incorporated into plasmids and transformed into *E. coli* strain SG13009 containing the pREP4 plasmid. The gene encoding PC₁₀P (with an N-terminal 6xHis tag) had already been prepared in the pQE-9 expression vector (Qiagen) by Shen et al.³⁰ (plasmid pQE9-PC₁₀P). The plasmid encoding PC₃₀P was prepared by first digesting pQE9-C₁₀ from Shen et al.³⁰ with *SpeI* and *NheI* to isolate the ~300 bp sequence encoding a C₁₀ unit and then digesting pQE9-PC₁₀P with *SpeI* and ligating the C₁₀ sequence into it. The digestion and ligation procedure was repeated a second time to form pQE9-PC₃₀P. Expression was performed in TB media and cultures were seeded at a 1:100 dilution with a saturated starter culture. Cells were grown at 37 °C until OD₆₀₀ reached 0.9–1.0 and then induced with 1.0 mM isopropyl β -D-1-thiogalactopyranoside (IPTG). Protein expression proceeded for 5 h, after which the cells were harvested and lysed in 8 M urea. Cells were frozen at –80 °C, thawed, and sonicated. Cell lysates were clarified by centrifugation, and the proteins of interest were isolated using Ni-NTA affinity chromatography. Yields of PC₁₀P and PC₃₀P were ~130 mg/L culture, and purity was confirmed by denatured state polyacrylamide gel electrophoresis (SDS-PAGE) and matrix-assisted laser desorption ionization mass spectrometry (MALDI-MS). Protein sequences and gels are included in the Supporting Information.

Gel samples were prepared by mixing lyophilized protein with 100 mM phosphate buffer with the pH adjusted to 7.6 for a total gel concentration of 7% w/v. Samples were held at 4 °C until the protein was fully hydrated, generally 2–4 h, then heated above the sol–gel transition by immersion in a boiling water bath for 60 s. The gels were allowed to reassociate under quiescent conditions at room temperature to erase any thermal and flow history. Rheological measurements were performed on a TA Instruments ARES-RFS strain-controlled rheometer with a 25 mm diameter cone and plate geometry. The outer edge of the sample was coated in oil in order to slow evaporation; the linear rheological properties of the materials were found to be stable for more than 48 h in this configuration. Linear rheological data were acquired at 10% strain amplitude (dynamic modulus, $G^*(\omega)$), and nonlinear rheological data were acquired using a fixed frequency (1 rad/s) that is in the elastic plateau region of $G^*(\omega)$. Time-dependent healing after cessation of large-amplitude oscillatory shear (LAOS) was probed using 0.1% strain amplitude at 1 rad/s frequency.

Flow visualization experiments were performed on PC₁₀P hydrogels using green fluorescent polystyrene microspheres (10 μ m diameter) as tracer particles at 0.1 vol %. The gel was loaded into a syringe and injected through a plastic capillary tube with 0.79 mm inner diameter using a syringe pump to control the nominal flow rate of the gel. Images were acquired on a Zeiss 510 inverted laser scanning confocal microscope according to the method of Fraser and co-workers.⁴² Line scans perpendicular to the flow direction were acquired in a plane through the center of the capillary at 1.93 ms/line for 2000 scans in order to produce images of the flow.

CHO K1 cells were grown to ~80% confluence on tissue culture polystyrene in cell growth media. Cell passages 16–20 were used for experiments. Cells were trypsinized, pelleted, and resuspended in fresh media. The cell suspension was then added to lyophilized PC₃₀P protein to make a total concentration of 7% w/v protein and gently stirred with the pipet tip. The mixture was placed on ice for 2 h to allow hydration, at which time a nearly homogeneous and transparent gel had formed. The gel was loaded into a 1 mL syringe and injected through a 22 gauge needle onto a coverslip. The gel was placed on ice for 15 min, at which point it was covered completely in cell growth media containing 2 μ M calcein and 4 μ M ethidium D (EthD) from the Invitrogen live/dead assay kit. The cells were held for 45 min before imaging on a Zeiss 510 inverted laser scanning confocal microscope. The green channel was excited at 488 nm, and the red channel was excited at 532 nm. DIC images were taken using 532 nm illumination.

Results and Discussion

The telechelic protein sequences PC_xP are constructed from two helical end blocks (denoted P) linked by x repeats of the nonapeptide sequence (C): PC₁₀P (20.5 kg/mol) has 10 nonapeptide repeats, while PC₃₀P (36.6 kg/mol) has 30 nonapeptide repeats. The coiled-coil P domains associate into pentameric bundles due to a combination of hydrophobic and ionic interactions, forming physical cross-links in the gels (Figure 1). Complete sequences of PC₁₀P and PC₃₀P are given in the Supporting Information.

While the linear rheology of the telechelic protein hydrogels is typical for physically gelling materials, the nonlinear rheology reveals strong shear-thinning behavior (Figure 2). In the linear regime, both PC₁₀P (Figure 2a) and PC₃₀P (Figure 2b) hydrogels exhibit plateaus in the elastic modulus, G' , for frequencies above 0.1 rad/s. At fixed mass concentration, the magnitude of the plateau modulus decreases with increasing length of the protein midblock from PC₁₀P to PC₃₀P—in contrast to the behavior observed earlier for tetrameric associative domains (AC_xA) examined by Shen et al.³⁴ The present materials approximately follow the expected scaling of the plateau modulus with the number of chains per unit volume: PC₃₀P has almost twice the molecular weight of PC₁₀P; hence, the 7% w/v gel of PC₃₀P has approximately half the number concentration of chains and has approximately half the modulus of the 7% w/v PC₁₀P gel. A stronger effect of midblock length is observed in the relaxation time of the gels, indicated by the frequency of the crossover of the elastic modulus, G' , and the viscous modulus, G'' ; relaxation is fast enough to enable observation of terminal behavior at low frequencies for PC₁₀P, but not for PC₃₀P.

Collectively, the behavior of the plateau modulus and the relaxation time suggests complex changes in the network structure. The proportionality between linear elastic modulus and chain density implies that the topology of the network (e.g., fraction of elastically effective chains) is similar for PC₁₀P and PC₃₀P gels at 7% w/v, in contrast to the AC_xA materials in which increasing the midblock length strongly promotes the formation of elastically effective linkages.³⁴ Thus, the choice of end block (P vs A) qualitatively changes the effects of increasing midblock length on the plateau modulus. The profound effect of the end block results from the discreteness of the aggregation number and the specificity of association (bias in favor of parallel vs antiparallel) of the coiled-coil proteins. While the moduli suggest similarities in the network structure between PC₁₀P and PC₃₀P, the increase in relaxation time from PC₁₀P to PC₃₀P suggests that the slowest relaxation process is qualitatively different in the two networks. This is in contrast to the expectation that the kinetics of a single coiled-coil exchange should be independent of the molecular weight of the midblock. Taken together, these results illustrate the interplay between end-block structure and midblock

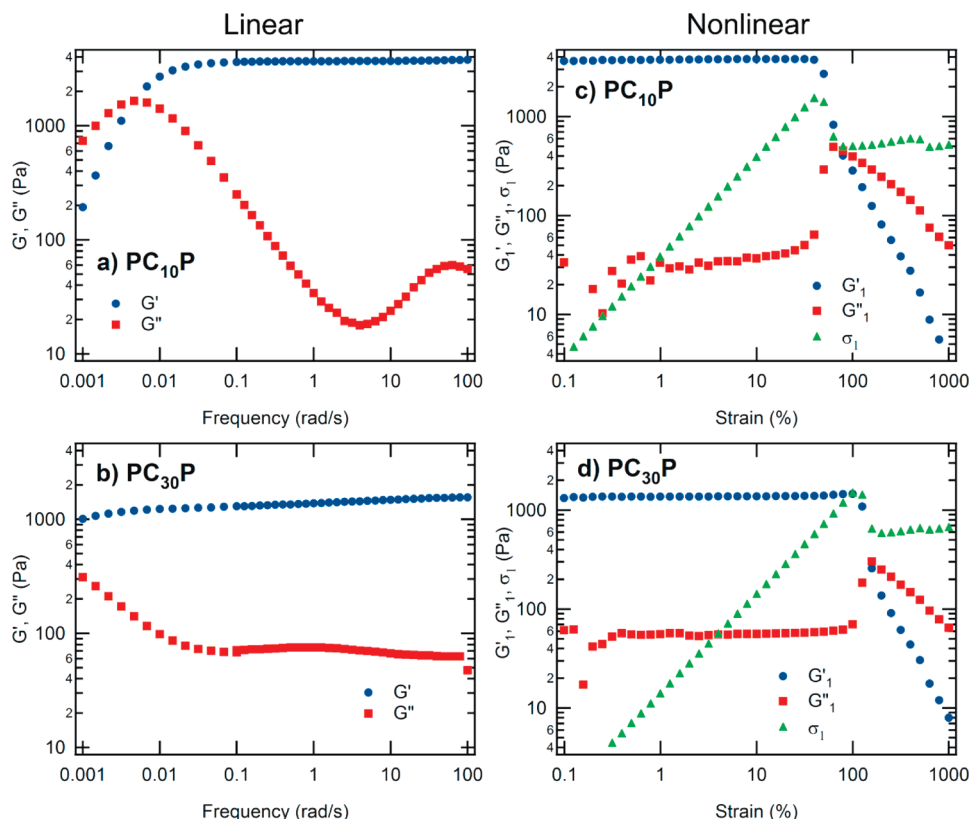


Figure 2. Linear and nonlinear rheology of PCP hydrogels. The linear dynamic moduli (left), elastic modulus, $G'(\omega)$, and loss modulus, $G''(\omega)$, and the nonlinear strain dependence of the first harmonic component of the stress (right), $\sigma_1 \equiv |\sigma_1^*(\gamma_0)|$, $G'_1(\gamma_0)$, and $G''_1(\gamma_0)$ at a fixed $\omega = 1$ rad/s, for PC₁₀P (top) and PC₃₀P (bottom) hydrogels that were prepared in 100 mM phosphate buffer, pH 7.6, at 7% w/v concentration. In the linear regime (left), each gel exhibits a plateau modulus at high frequency and a characteristic relaxation time that depends on the midblock length. At a fixed frequency of 1 rad/s, increasing the strain amplitude reveals a yield stress, above which the gels shear thin by 3 orders of magnitude.

length that govern the network structure and properties of artificial protein hydrogels.

Strain sweeps at 1 rad/s (Figure 2c,d) show that the linear regime extends to $\sim 50\%$ strain amplitude in both materials. Further increase in strain amplitude produces a sharp decrease in the amplitude of the stress response, which also becomes non-sinusoidal. Nevertheless, the first harmonic dominates throughout the range of strain amplitude explored ($\gamma_0 < 1000\%$). The abrupt decrease in the amplitude of the first harmonic, σ_1 , occurs over a narrow range of γ_0 (triangles, Figure 2, right), in which the character of the stress response changes from nearly elastic to predominantly dissipative: $G'_1(\gamma_0)$ decreases (circles) and G''_1 increases steeply (triangles) to meet G'_1 . The increase in energy dissipation at this transition correlates with disruption of non-covalent associations within the network. With further increase in γ_0 the stress amplitude σ_1 remains approximately constant. This stress plateau indicates an inverse relationship between strain amplitude and complex modulus; i.e., $|G^*|$ decreases as $1/\gamma_0$.⁴³ Since G''_1 dominates $|G^*|$, it scales approximately as $G''_1 \propto \gamma_0^{-1}$ (Figure 2, left), as observed previously by other authors.⁴³ (In the present materials, there is a slight increase of σ_1 with γ_0 , approximately $\sigma_1 \propto \gamma_0^{0.1}$, so $G''_1 \propto \gamma_0^{-0.9}$.) In the high- γ_0 plateau, G'_1 also falls as a power of γ_0 —with an exponent that is twice that for G''_1 (note slopes of -1.8 in G' and -0.9 in G'' in both parts c and d of Figure 2). The stress curves $\sigma_1(\gamma_0)$ of both materials are consistent with a yielding response with a characteristic yield stress ≈ 1400 Pa and plateau stress ≈ 550 Pa at high γ_0 . Notably, although the equilibrium modulus of the PC₁₀P gel is greater than that of PC₃₀P, the yield stresses and plateau stresses of the two materials are approximately the same; the origin of the common yield stress is not clear.

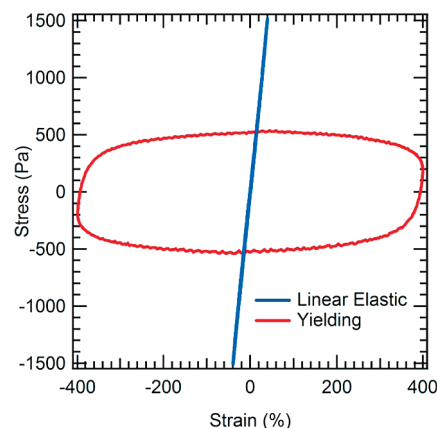


Figure 3. Lissajous figures for PC₁₀P. At low strain amplitude the material shows linear curves characteristic of an elastic solid, while at high strain a flattened ellipse is observed, indicative of yielding behavior. Curves are acquired at 1 rad/s for a 7% gel in 100 mM phosphate buffer at pH 7.6.

Characteristics of yielding behavior are also evident in the shape of the nonsinusoidal stress response at high strain: Lissajous figures recorded in large-amplitude oscillatory shear (LAOS) flow clearly become flattened (Figure 3) at high strain amplitudes, with the stress roughly constant with increasing strain during the oscillation, consistent with a yielding response.^{44,45} For comparison, at low strains the Lissajous figure is simply a line having a slope that reflects the linear viscoelastic properties of the gel at the test frequency.

The speed with which these telechelic protein hydrogels recover their full elastic strength is exceptional (Figure 4). After LAOS at strain amplitudes up to 500% (1 rad/s), PC₁₀P recovers to ~98% of its initial elastic modulus nearly instantaneously. While PC₃₀P

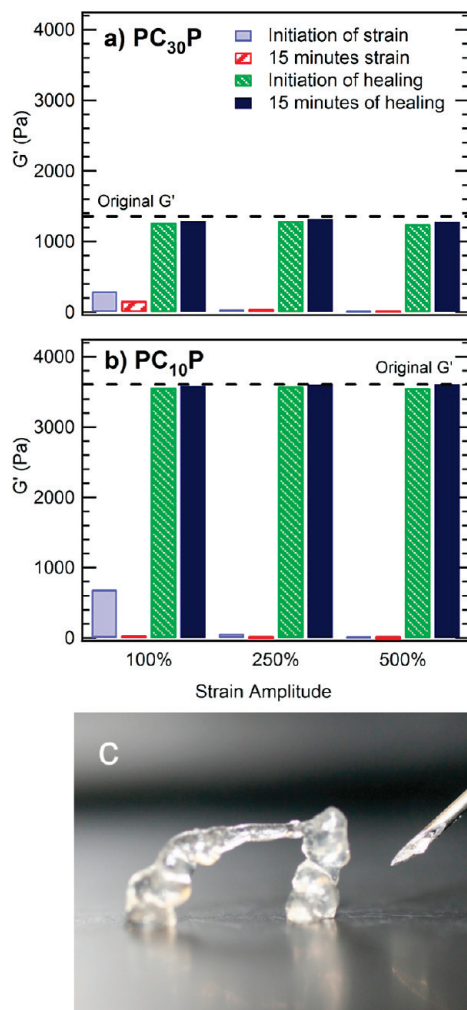


Figure 4. Rapid recovery of hydrogels after large-amplitude oscillatory strain (LAOS). (a) PC₃₀P recovers over 90% of its undisrupted modulus within seconds (the time required to record the first measurement at small strain, 0.1%) after the cessation of LAOS, independent of the imposed strain amplitude. (b) PC₁₀P recovers over 98% of its original modulus within the first measurement time, with greater than 99% recovery after 15 min. The measurement frequency was 1 rad/s for both samples. (c) The fast recovery of PC₁₀P allows self-supporting structures to be produced.

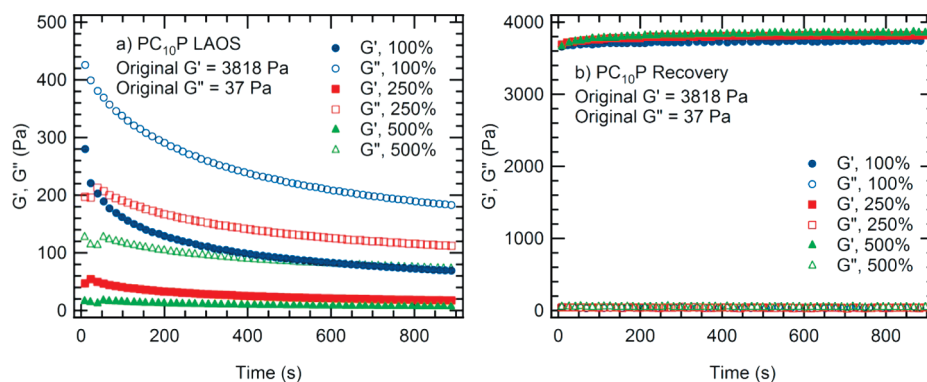


Figure 5. Time decay and recovery of mechanical properties for PC₁₀P under large-amplitude oscillatory shear. Measurements were performed at 1 rad/s. LAOS was performed for 15 min at the indicated strain amplitude, and recovery was conducted for 15 min at 0.1% strain amplitude.

shows a small permanent drop in the elastic modulus after the first LAOS strain cycle, the modulus does not decrease further after subsequent strain cycles of increasing amplitude. Although the yielding mechanism that results in shear thinning in these protein hydrogels may play a role in self-healing, comparison to other telechelic protein hydrogels and block copolymer hydrogels suggests that this rapid recovery is a property of the coiled-coil end blocks, since alternative end blocks lead to slower healing kinetics.⁴⁶ The time-dependent rheological response for PC₁₀P illustrates that, in addition to rapid recovery, the onset of shear thinning is rapid (Figure 5). The gels show a noticeable time decay of the modulus only for relatively low strain amplitudes near the end of the linear regime. The combination of extreme shear thinning with rapid recovery of mechanical strength allows these gels to be injected through narrow (22 gauge) needles with only gentle hand pressure and formed into self-supporting structures, such as the free-standing arch shown in Figure 4.

In order to investigate the nature of the yielding behavior, the velocity profile during capillary flow of a PC₁₀P hydrogel was visualized in a confocal microscope using fluorescent microspheres as flow markers.⁴² A series of line scans were acquired at a rate of 500 scans/s at a fixed axial position ~50 diameters from the capillary entrance. Particle velocities are visualized by composing the scan lines as time-sequential columns of pixels in an image (Figure 6, top), with position along the diameter of the capillary along the vertical axis and time along the horizontal axis. Particles that move more slowly intercept the line scan for a longer time and appear more elongated in the time dimension than fast particles, which appear compressed in the time direction. Each discernible particle in the image was measured and its size in pixels was converted to a velocity using the known diameter of the particles. Particles move with an average velocity of 0.151 mm/s, corresponding to a flow rate of 4.5 μ L/min, and the standard deviation of the velocity distribution is 15% of the mean. The average particle velocity does not vary with distance from the center of the capillary, indicating a plug flow velocity distribution. This suggests that gel in a solidlike state is flowing through the central region of the capillary due to yielding at or near the wall, which may be consistent with either wall slip or localized yielding near the wall within a narrow region of gel.

To distinguish whether shear banding or wall slip is the dominant mechanism for yielding, the gap thickness of the sample was varied in a parallel-plate geometry. These experiments show that the stress–strain curves are essentially independent of gap thickness, indicating that the yielding behavior is predominantly due to yielding within the bulk of the material (Supporting Information). Increased energy dissipation at the yield point also indicates yielding within the bulk of the gel, and a qualitatively similar rheological response has been observed in colloidal silica gels.⁴⁷ In the colloidal gels slip was minimized by the use of serrated

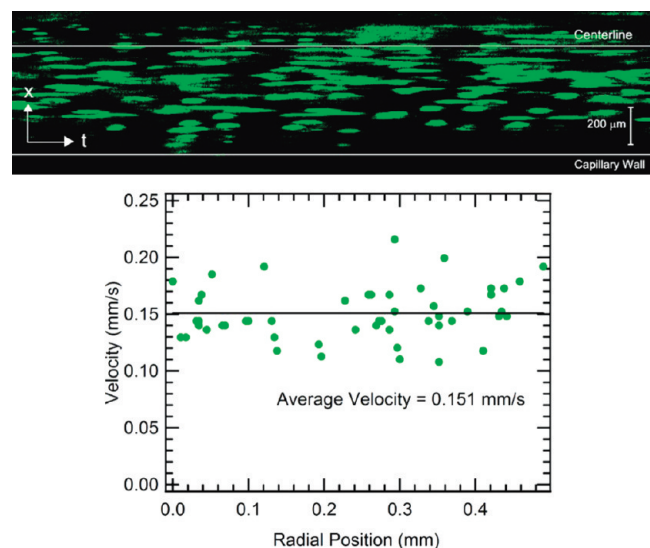


Figure 6. Visualization of the flow profile in capillary flow of a PC_{10}P hydrogel labeled with fluorescent microsphere tracers. Time-dependent line scans along the diameter of a capillary (1.93 ms/scan) show that the transit time for a particle to cross the plane of the line scan is approximately the same for all distances from the center of the capillary, indicating a plug flow type profile. This implies yielding in the gel layer adjacent to the capillary wall, consistent with the shear-banding mechanism.

plates, and the yielding behavior accompanied by the increase in G'' was attributed to network rupture. It is hypothesized that similar network rupture effects occur in the protein system due to dissociation of the coiled-coil protein end blocks under the high shear rates in a shear band. Given the density of chains in the system and the fact that the coiled-coil end blocks have interaction energies several times greater than the thermal energy, kT , the peak in G'' is small compared to the energy per unit volume of the coiled-coil linkages; therefore, dissociation of the coiled-coils within a small volume is consistent with the observed magnitude of the peak in G'' . The largest loss modulus is observed at the lowest strain amplitudes, consistent with the inability of chains to rapidly reassociate in high shear flows.

Yielding within the gel indicates the onset of shear banding,^{21,48,49} a phenomenon commonly observed in entangled polymer melts,^{50,51} DNA solutions,²³ wormlike micellar solutions,^{22,52–54} telechelic polymers,^{46,55} and micellar solutions cross-linked by telechelic polymers^{56,57} in both steady shear and LAOS. In the solid-like gel above the yield stress, strain is localized in regions where the cross-links or entanglements are depleted, allowing the intervening regions to move in a solidlike manner that maintains their cross-linked structure. The plateau in first harmonic stress results from yielding within this less entangled region. The overshoot in stress also suggests an activation barrier to shear thinning that is consistent with the nucleation hypothesis of shear band formation.⁵⁸ Because shear banding often results from a nonmonotonic stress–strain relationship, the stress in a homogeneous material may be increased into a metastable region, resulting in a stress drop after nucleation of a shear band.

Recently, the relationship between fracture and shear banding in transient networks has also been explored, and it has been shown that the microscopic fracture mechanism can be modeled with Griffith crack theory.⁵⁷ Crack growth in previously investigated hydrogels is visually observable, confirming the fracture mechanism. In these micellar solutions cross-linked by telechelic polymers, a plateau region has been observed even during unsteady shear banding.⁵⁶ Fracture was also observed in a telechelic protein using collagen helices as the associating groups investigated by Skrzyszewska et al.,⁴⁶ but in contrast to the transient micellar

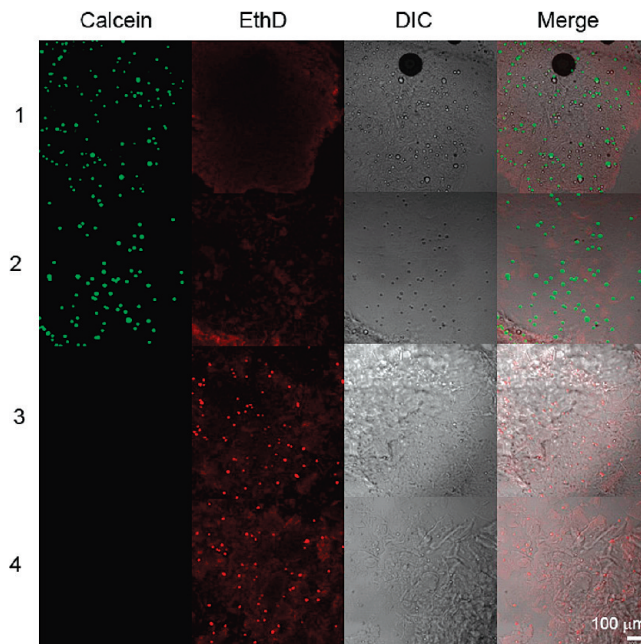


Figure 7. Cell survival after injection in a PC_{30}P hydrogel matrix. The live/dead cell assay indicates greater than 95% cell survival after the acute shear stress of injection. (1) Live cells without injection, (2) live cells after injection; (3) fixed cells without injection, and (4) fixed cells after injection. Calcein stain identifies live cells, while EthD stain identifies dead cells and adheres nonspecifically to the hydrogel.

networks, a plateau in stress was not observed with increasing strain rate with the collagen telechelic proteins. In the case of our coiled-coil protein hydrogels, visual observation during shear (using a transparent lower plate) shows no indication of fracture or void formation, and a plateau is observed after yielding. Sprakel et al.⁵⁹ show with particle-based simulations that a transition between shear banding and fracture may occur with increasing concentration; therefore, the absence of fracture in the coiled-coil protein gels considered here may indicate that the concentration of protein is below the transition from shear banding to fracture.

As noted earlier, shear banding is of interest in injectable gels for tissue engineering matrices because it may allow seeded cells to survive the injection process. To evaluate the ability of the present coiled-coil protein gel to protect cells during injection, CHO K1 cells were seeded into a PC_{30}P gel by allowing lyophilized protein to hydrate around the cells, and the cells were stained for viability after injection of the gel at a rate of $\sim 5 \text{ mL/min}$ (corresponding to a mean velocity of 0.43 m/s and nominal shear rate of 875 1/s). Cell survival was greater than 95% after injection and was not significantly different from that observed in the uninjected control (Figure 7). Negative controls with fixed cells show that the staining conditions are sufficient to identify dead cells. EthD dye non-specifically adhered to the protein gel, showing that the cells were encapsulated within the gel. Although relatively high shear stresses are required to induce shear thinning these gels, the shear-banding flow mechanism localizes the deformation to small regions of the gel. This allows a majority of seeded cells to experience extremely low shear strain and ensures high cell survival rates, in addition to allowing flow during injection.

Conclusions

A system of telechelic proteins with associating coiled-coil end blocks and a flexible polyelectrolyte midblock was used to formulate shear thinning injectable hydrogels. The materials demonstrate up to 3 orders of magnitude drop in their elastic modulus with increasing shear due to a sharp transition from a

linear elastic solid to yielding flow. Genetic engineering of the protein can be used to control the physical properties of the gel such as changing modulus with changing midblock length without affecting the yielding behavior responsible for shear thinning. Flow visualization and LAOS rheology were used to characterize the yielding behavior, demonstrating that it is consistent with the formation of shear bands within the hydrogel. The extreme shear thinning and nearly instantaneous recovery to full elastic strength after the cessation of shear make these materials easy to inject and to sculpt into self-supporting structures for application as tissue engineering matrices. The shear-banding flow mechanism localizes strain into small volumes of the gel, allowing very high survival of encapsulated cells.

Acknowledgment. This work was supported by the NSF Center for the Science and Engineering of Materials and NIH Grant EB1971. B.D.O. was supported by Award F32GM0834 from the National Institute of General Medical Sciences and by a Beckman Institute Postdoctoral Fellowship. We thank Professor S. E. Fraser for suggesting the method of flow visualization reported in Figure 3.

Supporting Information Available: Figures S1–S4. This material is available free of charge via the Internet at <http://pubs.acs.org>.

References and Notes

- (1) Yu, L.; Ding, J. D. *Chem. Soc. Rev.* **2008**, 37 (8), 1473–1481.
- (2) Mano, J. F. *Adv. Eng. Mater.* **2008**, 10 (6), 515–527.
- (3) Peppas, N. A.; Huang, Y.; Torres-Lugo, M.; Ward, J. H.; Zhang, J. *Annu. Rev. Biomed. Eng.* **2000**, 2, 9–29.
- (4) Sakiyama-Elbert, S. E.; Hubbell, J. A. *Annu. Rev. Mater. Res.* **2001**, 31, 183–201.
- (5) Engler, A. J.; Sen, S.; Sweeney, H. L.; Discher, D. E. *Cell* **2006**, 126 (4), 677–689.
- (6) Burdick, J. A.; Anseth, K. S. *Biomaterials* **2002**, 23 (22), 4315–4323.
- (7) Holland, T. A.; Tabata, Y.; Mikos, A. G. *J. Controlled Release* **2003**, 91 (3), 299–313.
- (8) Shu, X. Z.; Liu, Y. C.; Palumbo, F. S.; Lu, Y.; Prestwich, G. D. *Biomaterials* **2004**, 25 (7–8), 1339–1348.
- (9) Stile, R. A.; Burghardt, W. R.; Healy, K. E. *Macromolecules* **1999**, 32 (22), 7370–7379.
- (10) Bhattarai, N.; Ramay, H. R.; Gunn, J.; Matsen, F. A.; Zhang, M. Q. *J. Controlled Release* **2005**, 103 (3), 609–624.
- (11) Jeong, B.; Bae, Y. H.; Lee, D. S.; Kim, S. W. *Nature* **1997**, 388 (6645), 860–862.
- (12) Chiu, Y. L.; Chen, S. C.; Su, C. J.; Hsiao, C. W.; Chen, Y. M.; Chen, H. L.; Sung, H. W. *Biomaterials* **2009**, 30 (28), 4877–4888.
- (13) Ramachandran, S.; Tseng, Y.; Yu, Y. B. *Biomacromolecules* **2005**, 6 (3), 1316–1321.
- (14) Gupta, D.; Tator, C. H.; Shoichet, M. S. *Biomaterials* **2006**, 27 (11), 2370–2379.
- (15) Haines-Butterick, L.; Rajagopal, K.; Branco, M.; Salick, D.; Rughani, R.; Pilarz, M.; Lamm, M. S.; Pochan, D. J.; Schneider, J. P. *Proc. Natl. Acad. Sci. U.S.A.* **2007**, 104 (19), 7791–7796.
- (16) Liu, X. L.; Mao, Y. G.; Mathias, E. V.; Ma, C.; Franco, O.; Ba, Y.; Kornfield, J. A.; Wang, T. L.; Xue, L. J.; Zhou, B. S.; Yen, Y. *J. Sol-Gel Sci. Technol.* **2008**, 45 (3), 269–278.
- (17) Ni, X. P.; Cheng, A.; Li, J. *J. Biomed. Mater. Res., Part A* **2009**, 88A (4), 1031–1036.
- (18) Li, J.; Ni, X. P.; Leong, K. W. *J. Biomed. Mater. Res., Part A* **2003**, 65A (2), 196–202.
- (19) Liu, K. L.; Zhu, J. L.; Li, J. *Soft Matter* **2010**, 6 (10), 2300–2311.
- (20) Moller, P. C. F.; Rodts, S.; Michels, M. A. J.; Bonn, D. *Phys. Rev. E* **2008**, 77 (4), 041507.
- (21) Olmsted, P. D. *Rheol. Acta* **2008**, 47 (3), 283–300.
- (22) Walker, L. M. *Curr. Opin. Colloid Interface Sci.* **2001**, 6 (5–6), 451–456.
- (23) Boukany, P. E.; Hu, Y. T.; Wang, S. Q. *Macromolecules* **2008**, 41 (7), 2644–2650.
- (24) Zhou, L.; Ewoldt, R. H.; Cook, L. P.; McKinley, G. H. Probing shear-banding transitions of entangled liquids using Large Amplitude Oscillatory Shearing (LAOS) deformations. In *XVth Int. Congr. Rheol. - the Society of Rheology 80th Annual Meeting, Pts 1 and 2*; Co, A., Leal, L. G., Colby, R. H., Giacomini, A. J., Eds.; 2008; Vol. 1027, pp 189–191.
- (25) Ravindranath, S.; Wang, S. Q. *J. Rheol.* **2008**, 52 (2), 341–358.
- (26) Petka, W. A.; Harden, J. L.; McGrath, K. P.; Wirtz, D.; Tirrell, D. A. *Science* **1998**, 281 (5375), 389–392.
- (27) Krishna, O. D.; Kiick, K. L. *Biopolym. Peptide Sci.* **2010**, 94 (1), 32–48.
- (28) Xu, C. Y.; Breedveld, V.; Kopecek, J. *Biomacromolecules* **2005**, 6 (3), 1739–1749.
- (29) Xu, C. Y.; Kopecek, J. *Pharm. Res.* **2008**, 25 (3), 674–682.
- (30) Shen, W.; Zhang, K. C.; Kornfield, J. A.; Tirrell, D. A. *Nature Mater.* **2006**, 5 (2), 153–158.
- (31) Yang, J. Y.; Xu, C. Y.; Wang, C.; Kopecek, J. *Biomacromolecules* **2006**, 7 (4), 1187–1195.
- (32) Yang, J. Y.; Xu, C. Y.; Kopeckova, P.; Kopecek, J. *Macromol. Biosci.* **2006**, 6 (3), 201–209.
- (33) Wang, C.; Kopecek, J.; Stewart, R. J. *Biomacromolecules* **2001**, 2 (3), 912–920.
- (34) Shen, W.; Kornfield, J. A.; Tirrell, D. A. *Soft Matter* **2007**, 3 (1), 99–107.
- (35) Shen, W. Structure, Dynamics, and Properties of Artificial Protein Hydrogels Assembled Through Coiled-Coil Domains. PhD Thesis, California Institute of Technology, Pasadena, 2005.
- (36) Nowak, A. P.; Breedveld, V.; Pakstis, L.; Ozbas, B.; Pine, D. J.; Pochan, D.; Deming, T. J. *Nature* **2002**, 417 (6887), 424–428.
- (37) Wright, E. R.; McMillan, R. A.; Cooper, A.; Apkarian, R. P.; Conticello, V. P. *Adv. Funct. Mater.* **2002**, 12 (2), 149–154.
- (38) Nagapudi, K.; Brinkman, W. T.; Thomas, B. S.; Park, J. O.; Srinivasarao, M.; Wright, E.; Conticello, V. P.; Chaikof, E. L. *Biomaterials* **2005**, 26 (23), 4695–4706.
- (39) Nagapudi, K.; Brinkman, W. T.; Leisen, J.; Thomas, B. S.; Wright, E. R.; Haller, C.; Wu, X. Y.; Apkarian, R. P.; Conticello, V. P.; Chaikof, E. L. *Macromolecules* **2005**, 38 (2), 345–354.
- (40) Martens, A. A.; Portale, G.; Werten, M. W. T.; de Vries, R. J.; Eggink, G.; Stuart, M. A. C.; de Wolf, F. A. *Macromolecules* **2009**, 42 (4), 1002–1009.
- (41) Werten, M. W. T.; Teles, H.; Moers, A.; Wolbert, E. J. H.; Sprakel, J.; Eggink, G.; de Wolf, F. A. *Biomacromolecules* **2009**, 10 (5), 1106–1113.
- (42) Jones, E. A. V.; Baron, M. H.; Fraser, S. E.; Dickinson, M. E. *Am. J. Physiol.* **2004**, 287 (4), H1561–H1569.
- (43) Li, X.; Wang, S. Q.; Wang, X. R. *J. Rheol.* **2009**, 53 (5), 1255–1274.
- (44) Ewoldt, R. H.; Hosoi, A. E.; McKinley, G. H. An ontology for large amplitude oscillatory shear flow. In *XVth Int. Congr. Rheol. - the Society of Rheology 80th Annual Meeting, Pts 1 and 2*; Co, A., Leal, L. G., Colby, R. H., Giacomini, A. J., Eds.; 2008; Vol. 1027, pp 1135–1137.
- (45) Ewoldt, R. H.; Hosoi, A. E.; McKinley, G. H. *J. Rheol.* **2008**, 52 (6), 1427–1458.
- (46) Skrzyszewska, P. J.; Sprakel, J.; de Wolf, F. A.; Fokkink, R.; Stuart, M. A. C.; van der Gucht, J. *Macromolecules* **2010**, 43 (7), 3542–3548.
- (47) Walls, H. J.; Caines, S. B.; Sanchez, A. M.; Khan, S. A. *J. Rheol.* **2003**, 47 (4), 847–868.
- (48) Milner, S. T.; McLeish, T. C. B.; Likhtman, A. E. *J. Rheol.* **2001**, 45 (2), 539–563.
- (49) Adams, J. M.; Olmsted, P. D. *Phys. Rev. Lett.* **2009**, 102 (6), 067801.
- (50) Tapadia, P.; Ravindranath, S.; Wang, S. Q. *Phys. Rev. Lett.* **2006**, 96 (19), 196001.
- (51) Olmsted, P. D.; Radulescu, O.; Lu, C. Y. D. *J. Rheol.* **2000**, 44 (2), 257–275.
- (52) Salmon, J. B.; Colin, A.; Manneville, S.; Molino, F. *Phys. Rev. Lett.* **2003**, 90 (22), 228303.
- (53) Berret, J. F.; Roux, D. C.; Porte, G. *J. Phys. II* **1994**, 4 (8), 1261–1279.
- (54) Spenley, N. A.; Cates, M. E.; McLeish, T. C. B. *Phys. Rev. Lett.* **1993**, 71 (6), 939–942.
- (55) Sprakel, J.; Spruijt, E.; Stuart, M. A. C.; Besseling, N. A. M.; Lettinga, M. P.; van der Gucht, J. *Soft Matter* **2008**, 4 (8), 1696–1705.
- (56) Michel, E.; Appell, J.; Molino, F.; Kieffer, J.; Porte, G. *J. Rheol.* **2001**, 45 (6), 1465–1477.
- (57) Tabuteau, H.; Mora, S.; Porte, G.; Abkarian, M.; Liguore, C. *Phys. Rev. Lett.* **2009**, 102 (15), 155501.
- (58) Dhont, J. K. G.; Briels, W. J. *Rheol. Acta* **2008**, 47 (3), 257–281.
- (59) Sprakel, J.; Spruijt, E.; van der Gucht, J.; Padding, J. T.; Briels, W. J. *Soft Matter* **2009**, 5 (23), 4748–4756.